# On the osculator Lorentz spheres of timelike parallel $p_{i}$-equidistant ruled surfaces in the Minkowski 3 -space $R_{1}^{3}$ 

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#### Abstract

In this paper, we present radii and curvature axes of osculator Lorentz spheres of the timelike parallel $p_{i}$-equidistant ruled surfaces with a timelike base curve in the Minkowski 3 -space $R_{1}^{3}$ and give the arc lengths of indicatrix curves of timelike base curves of these surfaces.


## 1. Introduction

I. E. Valeontis [3] defined parallel $p$-equidistant ruled surfaces in $E^{3}$ and gave some results related to the striction curves of these surfaces.
M. Masal and N. Kuruoğlu [2] studied arc lengths, curvature radii, curvature axes, spherical involute and areas of real closed spherical indicatrix curves of base curves of parallel $p$-equidistant ruled surfaces in $E^{3}$. And also, M. Masal and N. Kuruoğlu [1] defined timelike parallel $p_{i}$-equidistant ruled surfaces with a timelike base curve in the Minkowski 3 -space and have studied dralls, the shape operators, Gaussian curvatures, mean curvatures, shape tensor, $q^{\text {th }}$ fundamental forms of these surfaces.

This paper is organized as follows. In Section 3 we have found radii and curvature axes of osculator Lorentz spheres of the timelike parallel $p_{i}$-equidistant ruled surfaces with a timelike base curve in the Minkowski 3 -space. And later in Section 4 we have given arc lengths of indicatrix curves of these surfaces.

## 2. Preliminaries

Let $\quad \alpha: I \rightarrow R_{1}^{3}, \alpha(t)=\left(\alpha_{1}(t), \alpha_{2}(t), \alpha_{3}(t)\right) \quad$ be a differentiable unit speed timelike curve in the Minkowski 3 -space, where $I$ is an open interval

[^0]in R containing the origin. Let $V_{1}$ be the tangent vector field of $\alpha, D$ be the Levi-Civita connection on $R_{1}^{3}$ and $D_{V_{1}} V_{1}$ be a spacelike vector. If $V_{1}$ moves along $\alpha$, then a timelike ruled surface $M$ which is given by the parametrization
$$
\varphi(t, v)=\alpha(t)+v V_{1}(t)
$$
is obtained. Let $\left\{V_{1}, V_{2}, V_{3}\right\}$ be an orthonormal frame field along $\alpha$ in $R_{1}^{3}$, where $V_{2}$ and $V_{3}$ are spacelike vectors. If $k_{1}$ and $k_{2}$ are the natural curvature and torsion of $\alpha(t)$, respectively, then for $\alpha$ the Frenet formulas are given by (see [4])
\[

$$
\begin{equation*}
V_{1}^{\prime}=k_{1} V_{2}, V_{2}^{\prime}=k_{1} V_{1}-k_{2} V_{3}, \quad V_{3}^{\prime}=k_{2} V_{2} \tag{2.1}
\end{equation*}
$$

\]

Using $\quad V_{1}=\alpha^{\prime} \quad$ and $\quad V_{2}=\frac{\alpha^{\prime \prime}}{\left\|\alpha^{\prime \prime}\right\|}$, we have $k_{1}=\left\|\alpha^{\prime \prime}\right\|>0$, where "' " means derivative with respect to time $t$ (see [1]).

Definition 2.1 ([1]). The planes corresponding to subspaces $S p\left\{V_{1}, V_{2}\right\}$, $S p\left\{V_{2}, V_{3}\right\}$ and $S p\left\{V_{3}, V_{1}\right\}$ along striction curves of timelike ruled surface $M$ are called asymptotic plane, polar plane and central plane, respectively.

Let us suppose that $\alpha^{*}=\alpha^{*}\left(t^{*}\right)$ is another differentiable timelike curve with arc-length and $\left\{V_{1}^{*}, V_{2}^{*}, V_{3}^{*}\right\}$ is the Frenet frame of this curve in three dimensional Minkowski space $R_{1}^{3}$. Hence, we define timelike ruled surface $M^{*}$ parametrically as follows:

$$
\varphi^{*}\left(t^{*}, v^{*}\right)=\alpha^{*}\left(t^{*}\right)+v^{*} V_{1}^{*}\left(t^{*}\right),\left(t^{*}, v^{*}\right) \in I \times \mathrm{R} .
$$

Definition 2.2 ([1]). Let $M$ and $M^{*}$ be two timelike ruled surfaces and let $p_{1}, p_{2}$ and $p_{3}$ be the distances between the polar planes, central planes and asymptotic planes, respectively. If the directions of $M$ and $M^{*}$ are parallel and the distances $p_{i}, \quad 1 \leqslant i \leqslant 3$, of $M$ and $M^{*}$ are constant, then the pair of ruled surfaces $M$ and $M^{*}$ is called timelike parallel $p_{i^{-}}$ equidistant ruled surfaces with a timelike base curve. If specifically $p_{i}=0$, then this pair of ruled surfaces is named as timelike parallel $p_{i}$-equivalent ruled surfaces with a timelike base curve, where the base curves of ruled surfaces $M$ and $M^{*}$ are of class $C^{2}$.

Therefore the pair of timelike parallel $p_{i}$-equidistant ruled surfaces are defined parametrically as

$$
\begin{aligned}
M: \varphi(t, v) & =\alpha(t)+v V_{1}(t), \quad(t, v) \in I \times \mathrm{R}, \\
M^{*}: \varphi^{*}\left(t^{*}, v^{*}\right) & =\alpha^{*}\left(t^{*}\right)+v^{*} V_{1}\left(t^{*}\right), \quad\left(t^{*}, v^{*}\right) \in I \times \mathrm{R},
\end{aligned}
$$

where $t$ and $t^{*}$ are the arc parameters of curves $\alpha$ and $\alpha^{*}$, respectively. Let the striction curve of $M$ be the base curve of $M$ and let $\alpha^{*}$ be a base curve
of $M^{*}$. In this case we can write

$$
\alpha^{*}=\alpha+p_{1} V_{1}+p_{2} V_{2}+p_{3} V_{3}
$$

where $p_{1}(t), p_{2}(t)$ and $p_{3}(t)$ are of class $C^{2}$ (see [1]).
Theorem 2.1 (see [1], Theorem 3.2 and Corollary 3.1). Let $M$ and $M^{*}$ be timelike parallel $p_{i}$-equidistant ruled surfaces.
i) The Frenet vectors of timelike parallel $p_{i}$-equidistant ruled surfaces $M$ and $M^{*}$ at $\alpha(t)$ and $\alpha^{*}\left(t^{*}\right)$ points are equivalent for $\frac{d t^{*}}{d t}>0$.
ii) There is a relation between the natural curvatures $k_{1}(t)$ and $k_{1}^{*}\left(t^{*}\right)$ of base curves and the torsions $k_{2}(t)$ and $k_{2}^{*}\left(t^{*}\right)$ of $M$ and $M^{*}$ as follows:

$$
k_{i}^{*}=k_{i} \frac{d t}{d t^{*}}, \quad 1 \leq i \leq 2
$$

## 3. On the osculator Lorentz spheres of timelike parallel $p_{i}$-equidistant ruled surfaces with a timelike base curve

In this section, we will investigate radii and curvature axes of osculator Lorentz spheres of timelike parallel $p_{i}$-equidistant ruled surfaces $M$ and $M^{*}$ with a timelike base curve.

We compute the locus of center of the osculator sphere $S_{1}^{2}$ which is the fourth order contact with the base curve $\alpha$ of $M$. Let us consider the function $f$ defined by

$$
\begin{aligned}
f: I & \rightarrow \mathrm{R} \\
t & \rightarrow f(t)=\langle\alpha(t)-a, \alpha(t)-a\rangle-R^{2},
\end{aligned}
$$

where $a$ and $R$ are the center and radius of $S_{1}^{2}$, respectively. Since $S_{1}^{2}$ is the fourth order contact with the curve $\alpha$, we can write

$$
f(t)=f^{\prime}(t)=f^{\prime \prime}(t)=f^{\prime \prime \prime}(t)=0
$$

From $f(t)=0$ we have

$$
\begin{equation*}
\langle\alpha(t)-a, \alpha(t)-a\rangle=R^{2} \tag{3.1}
\end{equation*}
$$

from $f^{\prime}(t)=0$ and $V_{1}(t)=\alpha^{\prime}(t)$ we get

$$
\begin{equation*}
\left\langle V_{1}(t), \alpha(t)-a\right\rangle=0, \tag{3.2}
\end{equation*}
$$

from $f^{\prime \prime}(t)=0$ and equation (2.1) we have

$$
\begin{equation*}
\left\langle V_{2}(t), \alpha(t)-a\right\rangle=\frac{1}{k_{1}(t)} . \tag{3.3}
\end{equation*}
$$

Furthermore, for the vector $\alpha(t)-a$, we can write

$$
\begin{equation*}
\alpha(t)-a=m_{1}(t) V_{1}(t)+m_{2}(t) V_{2}(t)+m_{3}(t) V_{3}(t), m_{i}(t) \in \mathrm{R}, \tag{3.4}
\end{equation*}
$$

where $\left\{V_{1}, V_{2}, V_{3}\right\}$ is the orthonormal frame field of $M$. From here we have

$$
\begin{align*}
& \left\langle\alpha(t)-a, V_{1}(t)\right\rangle=-m_{1}(t),\left\langle\alpha(t)-a, V_{2}(t)\right\rangle=m_{2}(t), \\
& \left\langle\alpha(t)-a, V_{3}(t)\right\rangle=m_{3}(t) . \tag{3.5}
\end{align*}
$$

From equations (3.2) and (3.3) we get

$$
\begin{equation*}
m_{1}(t)=0, \quad m_{2}(t)=\frac{1}{k_{1}(t)} \tag{3.6}
\end{equation*}
$$

From equations (3.1), (3.4) and (3.6) we obtain

$$
\begin{equation*}
R=\sqrt{m_{2}^{2}+m_{3}^{2}} \tag{3.7}
\end{equation*}
$$

or

$$
\begin{equation*}
m_{3}= \pm \sqrt{R^{2}-m_{2}^{2}} \tag{3.8}
\end{equation*}
$$

Using (3.4), for the center $a$ of $S_{1}^{2}$, we can write

$$
a=\alpha(t)-\frac{1}{k_{1}} V_{2}-\lambda V_{3}, \quad \lambda=m_{3}(t) \in \mathrm{R} .
$$

From $f^{\prime \prime \prime}(t)=0$ we have

$$
k_{1}^{\prime}\left\langle V_{2}(t), \alpha(t)-a\right\rangle+k_{1}\left\langle V_{2}^{\prime}(t), \alpha(t)-a\right\rangle+k_{1}\left\langle V_{2}(t), V_{1}(t)\right\rangle=0 .
$$

So, from (2.1), (3.5) and (3.6) we obtain

$$
\begin{equation*}
m_{3}=\frac{-k_{1}^{\prime}}{k_{1}^{2} k_{2}}=\frac{m_{2}^{\prime}}{k_{2}} . \tag{3.9}
\end{equation*}
$$

Similarly, we compute the locus of center of osculator sphere $S_{1}^{* 2}$ which is the fourth order contact with the timelike base curve $\alpha^{*}$ of $M^{*}$. Let us consider the function $f^{*}$ defined as

$$
\begin{aligned}
f^{*}: I & \rightarrow \mathrm{R} \\
\quad t^{*} & \rightarrow f^{*}\left(t^{*}\right)=\left\langle\alpha^{*}\left(t^{*}\right)-a^{*}, \alpha^{*}\left(t^{*}\right)-a^{*}\right\rangle-R^{* 2}
\end{aligned}
$$

where $a^{*}$ and $R^{*}$ are the center and the radius of $S_{1}^{* 2}$. In addition, for the vector $\alpha^{*}\left(t^{*}\right)-a^{*}$, we can write

$$
\alpha^{*}\left(t^{*}\right)-a^{*}=m_{1}^{*}\left(t^{*}\right) V_{1}^{*}\left(t^{*}\right)+m_{2}^{*}\left(t^{*}\right) V_{2}^{*}\left(t^{*}\right)+m_{3}^{*}\left(t^{*}\right) V_{3}^{*}\left(t^{*}\right), m_{i}^{*}\left(t^{*}\right) \in \mathrm{R},
$$

where $\left\{V_{1}^{*}, V_{2}^{*}, V_{3}^{*}\right\}$ is the orthonormal frame field of $M^{*}$.
In a similar way $m_{1}^{*}\left(t^{*}\right), m_{2}^{*}\left(t^{*}\right), m_{3}^{*}\left(t^{*}\right), R^{*}$ and $a^{*}$ of $S_{1}^{* 2}$ are found to be

$$
\begin{equation*}
m_{1}^{*}\left(t^{*}\right)=0, m_{2}^{*}\left(t^{*}\right)=\frac{1}{k_{1}^{*}\left(t^{*}\right)}, m_{3}^{*}\left(t^{*}\right)=\frac{m_{2}^{*^{\prime}}}{k_{2}^{*}\left(t^{*}\right)}, R^{*}=\sqrt{m_{2}^{* 2}+m_{3}^{* 2}} \tag{3.10}
\end{equation*}
$$

and

$$
a^{*}=\alpha^{*}\left(t^{*}\right)-\frac{1}{k_{1}^{*}} V_{2}^{*}-\lambda^{*} V_{3}^{*}, \quad \lambda^{*}=m_{3}^{*}\left(t^{*}\right) \in \mathrm{R} .
$$

Now, we can compute the relations between the radii of osculator Lorentz spheres and curvature axes of the base curves of $M$ and $M^{*}$. From Theorem 2.1 ii), equations (3.6) and (3.10), we have

$$
\begin{equation*}
m_{1}^{*}\left(t^{*}\right)=m_{1}(t)=0, m_{2}^{*}\left(t^{*}\right)=\frac{d t^{*}}{d t} m_{2}(t) \tag{3.11}
\end{equation*}
$$

If $\frac{d t}{d t^{*}}$ is constant, then from Theorem 2.1 ii) we obtain

$$
\begin{equation*}
k_{1}^{*^{\prime}}=k_{1}^{\prime}\left(\frac{d t}{d t^{*}}\right)^{2} \tag{3.12}
\end{equation*}
$$

Hence, using (3.10), (3.12), (3.9) and Theorem 2.1 ii), we find

$$
\begin{equation*}
m_{3}^{*}=\frac{d t^{*}}{d t} m_{3} \tag{3.13}
\end{equation*}
$$

Combining (3.11), (3.13) and Theorem 2.1 i), we get

$$
\alpha^{*}-a^{*}=\frac{d t^{*}}{d t}(\alpha-a)
$$

Similarly, combining (3.7), (3.8), (3.11) and (3.13), we have

$$
R^{* 2}=\left(\frac{d t^{*}}{d t}\right)^{2} R^{2}
$$

or

$$
R^{*}=\left|\frac{d t^{*}}{d t}\right| R
$$

So, we have proved the following theorem.
Theorem 3.1. Let $M$ and $M^{*}$ be the timelike parallel $p_{i}$-equidistant ruled surfaces with a timelike base curve.
i) If $q_{\alpha}$ and $q_{\alpha^{*}}$ are the curvature axes (the locus of center of osculator Lorentz spheres) of the base curves $\alpha$ and $\alpha^{*}$ of $M$ and $M^{*}$, then we have

$$
q_{\alpha^{*}}-\alpha^{*}=\frac{d t^{*}}{d t}\left(q_{\alpha}-\alpha\right)
$$

ii) If $R$ and $R^{*}$ are the radiuses of osculator Lorentz spheres of base curves $\alpha$ and $\alpha^{*}$ of $M$ and $M^{*}$, then we have

$$
R^{*}=\left|\frac{d t^{*}}{d t}\right| R
$$

## 4. Arc lengths of indicatrix curves of the timelike parallel $p_{i}$-equidistant ruled surfaces with a timelike base curve

In this section, we will investigate arc lengths of indicatrix curves of timelike base curves of the timelike parallel $p_{i}$-equidistant ruled surfaces $M$ and $M^{*}$ with timelike base curve.

Since $V_{2}$ and $V_{3}$ are spacelike vectors, the curves $\left(V_{2}\right)$ and $\left(V_{3}\right)$ generated by the spacelike vectors $V_{2}$ and $V_{3}$ on the pseudosphere $S_{1}^{2}$ are called the pseudo-spherical indicatrix curves. The curve $\left(V_{1}\right)$ generated by the vector $V_{1}$ on the pseudohyperbolic space $H_{1}^{2}$ is called indicatrix curve. Let $S_{V_{i}}$ and $S_{V_{i}^{*}}$ denote the arc lengths of indicatrix curves $\left(V_{i}\right)$ and $\left(V_{i}^{*}\right)$ generated by the vector fields $V_{i}$ and $V_{i}^{*}$, respectively. So we can write

$$
S_{V_{i}}=\int\left\|V_{i}^{\prime}\right\| d t \text { and } S_{V_{i}^{*}}=\int\left\|V_{i}^{*^{\prime}}\right\| d t^{*}, 1 \leq i \leq 3
$$

Using the Frenet formulas and Theorem 2.1 ii), we get
$S_{V_{\mathrm{i}}^{*}}=\int k_{1} d t=S_{V_{i}}, \quad S_{V_{2}^{*}}=\int \sqrt{\left|k_{2}^{2}-k_{1}^{2}\right|} d t=S_{V_{2}}, S_{V_{3}^{*}}=\int\left|k_{2}\right| d t=S_{V_{3}}$, where $\frac{d t}{d t^{*}}>0$.

Similarly, for the arc lengths $S_{\alpha}$ and $S_{\alpha^{*}}$ of the indicatrix curves $(\alpha)$ and $\left(\alpha^{*}\right)$ generated by the timelike curves $\alpha$ and $\alpha^{*}$ on the pseudosphere $S_{1}^{2}$, respectively, we can write

$$
S_{\alpha}=\int\left\|\alpha^{\prime}\right\| d t=\int d t \text { and } S_{\alpha^{*}}=\int\left\|\alpha^{*^{\prime}}\right\| d t^{*}=\int d t^{*}
$$

If $\frac{k_{1}}{k_{1}^{*}}$ is constant, then using Theorem 2.1 ii), we obtain

$$
S_{\alpha^{*}}=\frac{k_{1}}{k_{1}^{*}} S_{\alpha}
$$

Thus we have proved the following theorems.
Theorem 4.1. If $S_{V_{i}}$ and $S_{V_{i}^{*}}, 1 \leq i \leq 3$, are the arc lengths of indicatrix curves of Frenet vectors $V_{i}$ and $V_{i}^{*}$ of timelike base curves $\alpha$ and $\alpha^{*}$ of the timelike parallel $p_{i}$-equidistant ruled surfaces $M$ and $M^{*}$, respectively, then we have

$$
S_{V_{i}^{*}}=S_{V_{i}}, \quad 1 \leq i \leq 3
$$

Theorem 4.2. Let $S_{\alpha}$ and $S_{\alpha^{*}}$ be the arc lengths of indicatrix curves of timelike base curves $\alpha$ and $\alpha^{*}$ of the timelike parallel $p_{i}$-equidistant ruled
surfaces $M$ and $M^{*}$, respectively. If $\frac{k_{1}}{k_{1}^{*}}$ is constant, then we have $S_{\alpha^{*}}=$ $\frac{k_{1}}{k_{1}^{*}} S_{\alpha}$.

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