

Another generalization of the bivariate FGM distribution with two-dimensional extensions

CARLES M. CUADRAS AND WALTER DÍAZ

ABSTRACT. The Farlie–Gumbel–Morgenstern family of bivariate distributions with given marginals is frequently used in theory and applications and has been generalized in many ways. With the help of two auxiliary distributions, we propose another generalization and study its properties. After defining the dimension of a distribution as the cardinal of the set of canonical correlations, we prove that some well-known distributions are practically two-dimensional. Then we introduce an extended FGM family in two dimensions and study how to approximate any distribution to this family.

1. Introduction

Constructing dependence models by means of copulas has interest in statistics, probability, econometrics, informatics, insurance, biometry, physics, etc. A copula function is a bivariate cumulative distribution function (cdf) with uniform $[0, 1]$ marginals, that captures the dependence properties of two random variables (r.v., in short). Many copulas and bivariate families of distributions have been described in Hutchinson and Lai (1991), Joe (1997), Drouet-Mari and Kotz (2001) and Nelsen (2006). Among others, the so-called Farlie–Gumbel–Morgenstern (FGM) bivariate family is frequently used in theory and applications. This motivated Huang and Kotz (1999), Amblard and Girard (2002), Nelsen *et al.* (1997), Rodríguez-Lallena and Úbeda-Flores (2004), Durante (2007), Cuadras (2009) and Cuadras *et al.* (2000) to study proper extensions.

Throughout this paper, x, y in $H(x, y)$, $F(x)$, $G(y)$, as well as u, v in $C(u, v)$, will be suppressed, unless it is strictly necessary. We write $H \in$

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$\mathcal{F}(F, G)$, where $\mathcal{F}(F, G)$ is the Fréchet family of distribution functions with marginals F, G . A member of $\mathcal{F}(F, G)$ is the FGM family

$$H_\theta = FG[1 + \theta(1 - F)(1 - G)], \quad -1 \leq \theta \leq 1.$$

We show that this family is one-dimensional and propose a two-dimensional extension in a geometric sense.

As in Nelsen *et al.* (1997), who obtained second-order approximations, the goal of this paper is to prove that some well-known distributions can be approximated by the extended family proposed in Section 3.

2. Geometric dimensionality

Suppose that the Radon–Nikodým derivative $dH/dFdG$ exists. A global measure of dependence is *Pearson contingency coefficient* ϕ^2 defined by

$$\phi^2 = \int_a^b \int_c^d \left(\frac{dH}{dFdG} - 1 \right)^2 dFdG.$$

We have $\phi^2 \geq 0$ and $\phi^2 = 0$ if and only if there is stochastic independence between X and Y .

If ϕ^2 is finite, then there exists the diagonal expansion (Lancaster, 1958)

$$dH = dFdG + \sum_{n \geq 1} \rho_n A_n B_n dFdG, \quad (1)$$

where A_n, B_n are unitary functions in $L^2([a, b])$ and $L^2([c, d])$ on F and G respectively, in the sense that $E[A_n(X)] = E[B_n(Y)] = 0$ and $E[A_n^2(X)] = E[B_n^2(Y)] = 1$. Then $A_n(X)$ and $B_n(Y)$ are the canonical variables. The sequence of canonical correlations is $\rho_1 \geq \rho_2 \geq \dots \geq 0$. It can be proved from expansion (1) that Pearson contingency coefficient can be expressed in terms of this sequence:

$$\phi^2 = \sum_{n \geq 1} \rho_n^2.$$

The first canonical correlation $\rho_1 = \sup \text{corr}(A(X), B(Y))$, where corr means correlation coefficient, is the maximal correlation between a function of X and a function of Y . The sequence of canonical correlations captures the full dependence between X and Y and the coefficient ϕ^2 is an overall measure of dependence, often presented as the ratio $\phi^2/(1 + \phi^2)$.

Definition 1. The geometric dimension of $H \in \mathcal{F}(F, G)$ such that the canonical expansion (1) exists, is the cardinal of the set (ρ_n) .

The cdf H with dimension $N \leq \infty$, can be approximated by another cdf H_D of smaller dimension $D < N$, where

$$dH_D = dFdG + \sum_{n=1}^D \rho_n A_n B_n dFdG.$$

The “proportion of dependence” of H accounted by H_D is:

$$P_D = \frac{\sum_{n=1}^D \rho_n^2}{\phi^2}.$$

This proportion is common in some methods of multivariate analysis. For instance, it is used in correspondence analysis to measure the quality of the graphical representation of a contingency table with respect to the chi-square distance, see Greenacre (1984) and Cuadras *et al.* (2000).

On the other hand, from the canonical expansion (1), we can also consider the following nested family

$$dH_\lambda = dFdG + \sum_{i=1}^k \lambda_i A_i B_i dFdG,$$

depending on the parameters $\lambda_1, \lambda_2, \dots, \lambda_k$, where $|\lambda_i| \leq \rho_i$, for $i = 1, \dots, k$.

Next, let us consider the FGM copula. Recall that a bivariate copula is a function $C : I^2 \rightarrow I$, with $I = [0, 1]$. such that $C(u, 0) = C(0, v) = 0$, $C(u, 1) = u, C(1, v) = v$, and for $0 \leq u_1 \leq u_2 \leq 1$ and $0 \leq v_1 \leq v_2 \leq 1$ satisfies $C(u_2, v_2) - C(u_2, v_1) - C(u_1, v_2) + C(u_1, v_1) \geq 0$.

Copulas are important because of Sklar’s theorem. Let $H(x, y)$ be the bivariate cdf of (X, Y) , with univariate marginal distribution functions $F(x)$, $G(y)$ and supports $[a, b], [c, d]$, respectively. Then this bivariate cdf can be expressed as $H(x, y) = C(F(x), G(y))$, where C is a copula related to H . Thus modeling copulas is an interesting task.

The related copula for the FGM distribution is

$$C_\theta = uv[1 + \theta(1 - u)(1 - v)], \quad -1 \leq \theta \leq 1.$$

The density is $c(u, v) = 1 + \theta(1 - 2u)(1 - 2v)$. Since the only canonical correlation is $\rho_1 = |\theta|/3$, clearly this distribution has geometric dimension one. In general, the dimension is higher.

We are interested in two-dimensional distributions as possible approximations to a given distribution. For example, for the following copulas (see Nelsen, 2006, Çelebioğlu, 1997, Cuadras, 2009):

$$\begin{aligned} \text{Ali–Mikhail–Haq:} & \quad uv/[1 - \theta(1 - u)(1 - v)], \quad -1 \leq \theta \leq 1, \\ \text{Gumbel–Barnett:} & \quad uv \exp(-\theta \ln u \ln v), \quad 0 \leq \theta \leq 1, \\ \text{Çelebioğlu–Cuadras:} & \quad uv \exp[\theta(1 - u)(1 - v)], \quad -1 \leq \theta \leq 1. \end{aligned}$$

The classical FGM family is the first order approximation from a Taylor expansion. However, the full dimension in these three copulas is countable (i.e., \aleph_0), whereas the approximation using only two dimensions is very accurate, as the proportion of dependence is $P_2 \simeq 0.99$.

The extension of the FGM family proposed below is two-dimensional in the above geometric sense.

3. Two-dimensional extension

Recall that the supports of F, G are $[a, b], [c, d]$. Let Φ, Ψ be univariate distribution functions with the same supports. An extension of FGM is the bivariate family

$$\begin{aligned} H(x, y) &= F(x)G(y) + \lambda_1[F(x) - \Phi(x)][G(y) - \Psi(y)] \\ &\quad + \lambda_2[\frac{1}{2}F^2(x) + (F_\Phi - \frac{1}{2})F(x) - F_\Phi(x)][\frac{1}{2}G^2(y) + (G_\Psi - \frac{1}{2})G(y) - G_\Psi(y)], \end{aligned} \quad (2)$$

where

$$F_\Phi(x) = \int_a^x \Phi(t)dF(t), \quad G_\Psi(y) = \int_c^y \Psi(t)dG(t), \quad F_\Phi = F_\Phi(b), \quad G_\Psi = G_\Psi(d).$$

The density function (Lebesgue measure) is

$$h = fg + \lambda_1 f(1 - \varphi f^{-1})g(1 - \psi g^{-1}) + \lambda_2 f(F - \Phi - \gamma)g(G - \Psi - \delta), \quad (3)$$

where

$$\gamma = \int_a^b (F - \Phi)dF = 1/2 - F_\Phi, \quad \delta = \int_c^d (G - \Psi)dG = 1/2 - G_\Psi.$$

Of course, (2) reduces to FGM for $\Phi = F^2, \Psi = G^2$ and $\lambda_2 = 0$.

This family is diagonal and the canonical correlations are

$$\begin{aligned} \rho_1 &= \lambda_1 \sqrt{(\alpha - 1)(\beta - 1)}, \\ \rho_2 &= \lambda_2 \sqrt{\frac{1/3 + F_{\Phi^2} - (1/2 - F_\Phi)^2 - 2 \int_a^b F\Phi dF}{1/3 + G_{\Psi^2} - (1/2 - G_\Psi)^2 - 2 \int_c^d G\Psi dG}}, \end{aligned}$$

where

$$\alpha = \int_a^b \left(\frac{d\Phi}{dF}\right)^2 dF, \quad \beta = \int_c^d \left(\frac{d\Psi}{dG}\right)^2 dG,$$

and $F_{\Phi^2} = \int_a^b \Phi^2 dF, G_{\Psi^2} = \int_c^d \Psi^2 dG$.

To prove this write (2) as

$$dH = dFdG + \lambda_1 a_1 dF b_1 dG + \lambda_2 a_2 dF b_2 dG,$$

where $a_1 = 1 - d\Phi/dF, b_1 = 1 - d\Psi/dG, a_2 = (F - \Phi - \gamma)$ and $b_2 = (G - \Psi - \delta)$.

It is readily proved that $E(a_1) = E(a_2) = 0$ and $E(b_1) = E(b_2) = 0$.

Note that

$$\begin{aligned} \int_a^b (1 - d\Phi/dF)(F - \Phi - \gamma)dF &= 0, \\ \int_a^b (F - \Phi - \gamma)^2 dF &= 1/3 + F_{\Phi^2} - (1/2 - F_\Phi)^2 - 2 \int_a^b F\Phi dF, \\ \int_c^d (G - \Psi - \delta)^2 dG &= 1/3 + G_{\Psi^2} - (1/2 - G_\Psi)^2 - 2 \int_c^d G\Psi dG. \end{aligned}$$

Hence $E(a_1 a_2) = \int_a^b a_1 a_2 dF = 0$ and similarly $E(b_1 b_2) = 0$.

The other covariances are:

$$\begin{aligned} \text{cov}(a_1, b_1) &= \int_a^b \int_c^d a_1 b_1 (dH - dF dG) \\ &= \lambda_1 \int_a^b a_1^2 dF \int_c^d b_1^2 dG + \lambda_2 \int_a^b a_1 a_2 dF \int_c^d b_1 b_2 dG \\ &= \lambda_1 (\alpha - 1) (\beta - 1), \\ \text{cov}(a_1, b_2) &= \lambda_1 \int_a^b a_1^2 dF \int_c^d b_1 b_2 dG + \lambda_2 \int_a^b a_1 a_2 dF \int_c^d b_2^2 dG \\ &= 0, \end{aligned}$$

and similarly $\text{cov}(a_2, b_1) = 0$. Also

$$\begin{aligned} \text{cov}(a_2, b_2) &= 0 + \lambda_2 \int_a^b a_2^2 dF \int_c^d b_2^2 dG \\ &= \lambda_2 [1/3 + F_{\Phi^2} - (1/2 - F_{\Phi})^2 - 2 \int_a^b F \Phi dF] \\ &\quad \times [1/3 + G_{\Psi^2} - (1/2 - G_{\Psi})^2 - 2 \int_c^d G \Psi dG]. \end{aligned}$$

The variances are

$$\begin{aligned} E(a_1^2) &= \alpha - 1 \\ E(a_2^2) &= 1/3 + F_{\Phi^2} - (1/2 - F_{\Phi})^2 - 2 \int_a^b F \Phi dF, \text{ etc.} \end{aligned}$$

Next, we study the corresponding copula. First note that $Q = \Phi \circ F^{-1}$ and $R = \Psi \circ G^{-1}$ are distribution functions with support in $[0, 1]$.

Since $F_{\Phi} = \int_a^b \Phi dF = \int_0^1 Q(t) dt = 1 - \mu_Q$, where μ_Q is the mean of the r.v. with cdf Q , and similarly μ_R , the copula corresponding to (2) is

$$\begin{aligned} C(u, v) &= uv + \lambda_1 [u - Q(u)][v - R(v)] \\ &\quad + \lambda_2 \left[\frac{1}{2} u^2 + \left(\frac{1}{2} - \mu_Q \right) u - \int_0^u Q(t) dt \right] \left[\frac{1}{2} v^2 \right. \\ &\quad \left. + \left(\frac{1}{2} - \mu_R \right) v - \int_0^v R(t) dt \right]. \end{aligned} \quad (4)$$

For $\Phi = F^k, \Psi = G^l$, family (2) has the copula

$$\begin{aligned} C(u, v) &= uv + \lambda_1 u(1 - u^{k-1})v(1 - v^{l-1}) \\ &\quad + \lambda_2 \left[\frac{1}{2} u^2 - u(k-1)/(2(k+1)) - u^{k+1}/(k+1) \right] \text{ [similar term in } v, l] \end{aligned}$$

and Spearman's correlation $\rho_S = 12 \int_{I^2} C dudv - 3$ is

$$\rho_S = \frac{(k-1)(l-1)}{(k+1)(l+1)} \left[3\lambda_1 + \frac{\lambda_2(k-2)(l-2)}{12(k+2)(l+2)} \right].$$

In particular, for $\Phi = F^2, \Psi = G^2$, the copula is

$$C = uv + uv(1-u)(1-v)[\lambda_1 + \bar{\lambda}_2(1-2u)(1-2v)], \quad (5)$$

where $\bar{\lambda}_2 = \lambda_2/36$.

Nelsen *et al.* (1997) studied the copula (5) as a member of the family of symmetric copulas with cubic sections in both u and v , providing Spearman's rho, Kendall's tau, certain dependence properties and showing that they are second-degree Maclaurin approximations to members of the Frank and Plackett copula families.

4. Approximating a bivariate cdf for another one

Let us consider $H_a \in \mathcal{F}(F, G)$ and suppose that the canonical expansion $dH_a = dFdG + \sum_{n \geq 1} \rho_n A_n dF B_n dG$ exists, where A_n, B_n are unitary canonical functions. Let $H_t \in \mathcal{F}(F, G)$ be the “true” cdf of two observable random variables (X, Y) . We are interested in approximating H_t by means of a finite linear combination of canonical functions obtained from H_a :

$$dH_t \simeq dFdG + \sum_{i=1}^k \lambda_i A_i dF B_i dG.$$

In a more precise way, we seek for the approximation

$$\frac{dH_t}{dFdG} \simeq 1 + \sum_{i=1}^k \lambda_i A_i B_i,$$

where $\lambda_1, \dots, \lambda_k$ are real coefficients such that

$$\int_a^b \int_c^d \left(\frac{dH_t - dFdG}{dFdG} - \sum_{i=1}^k \lambda_i A_i B_i \right)^2 dFdG \quad (6)$$

is minimized. If the densities h_t, f, g exist, then h_t is approximated by

$$\hat{h}_t = fg \left(1 + \sum_{i=1}^k \lambda_i A_i B_i \right).$$

Theorem 1. *Suppose $(X, Y) \sim H_t \in \mathcal{F}(F, G)$. The coefficients minimizing (6) are $\lambda_i = \rho_i$, where*

$$\rho_i = \text{corr}(A_i(X), B_i(Y)), \quad i = 1, \dots, k.$$

Proof. Write $z = (dH_t - dFdG)/dFdG$. Then

$$\begin{aligned} \int_a^b \int_c^d (z - \sum_{i=1}^k \lambda_i A_i B_i)^2 dFdG &= \phi_t^2 + \int_a^b \int_c^d \sum_{i=1}^k \lambda_i^2 A_i^2 B_i^2 dFdG \\ &\quad - 2 \int_a^b \int_c^d (\sum_{i=1}^k \lambda_i A_i B_i) (dH_t - dFdG) \\ &\quad + \sum_{i \neq j=1}^k \lambda_i \lambda_j \int_a^b A_i dF \int_c^d B_j dG \\ &= \phi_t^2 + \sum_{i=1}^k \lambda_i^2 - 2 \sum_{i=1}^k \lambda_i \rho_i, \end{aligned}$$

where ϕ_t^2 is the Pearson contingency coefficient. Taking the partial derivative with respect to λ_i on the left hand side of this equation, and equating to zero, we obtain $\lambda_i = \rho_i$, $i = 1, \dots, k$. \square

Note that each ρ_i is the correlation between the canonical variables A_i, B_i obtained from H_a , but the correlation is taken with respect to the “true” cdf H_t . That is, the observed data used in computing ρ_i comes from $(X, Y) \sim H_t$. This result is useful when the canonical functions of H_a are known, as it occurs in the above generalized FGM.

5. Examples

5.1. Ali–Mikhail–Haq copula. We first study the approximation of the Ali–Mikhail–Haq (AMH) copula (Nelsen, 2006)

$$C_t = \frac{uv}{1 - \theta(1 - u)(1 - v)}, \quad -1 \leq \theta \leq 1,$$

to the copula (5). The density for this generalized FGM copula is

$$c = 1 + \frac{\lambda_1}{3}\sqrt{3}(1 - 2u)\sqrt{3}(1 - 2v) + \frac{\lambda_2}{5}\sqrt{5}(6u^2 - 6u + 1)\sqrt{5}(6v^2 - 6v + 1).$$

The canonical functions are $A_1 = \sqrt{3}(1 - 2u)$, $A_2 = \sqrt{5}(6u^2 - 6u + 1)$ and similarly B_1, B_2 . To be sure that c is a density, the canonical correlations should belong to the region $\mathcal{R} = \{(\rho_1, \rho_2) | c \geq 0\}$, see Figure 1.

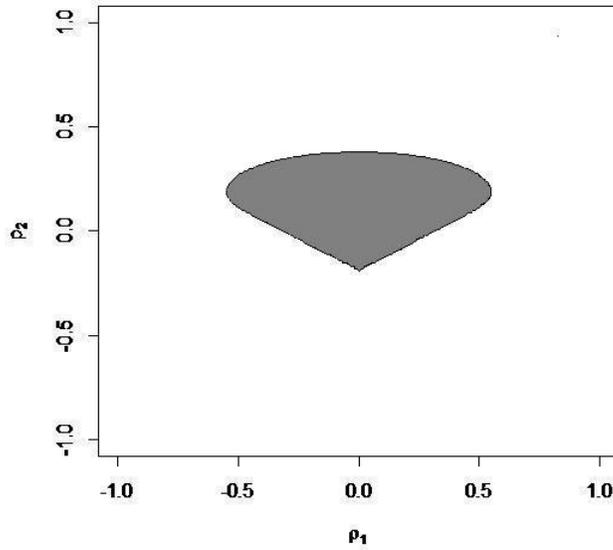


FIGURE 1. Region of the correlations (parameters) for which the density is positive.

We should calculate the correlations

$$\rho_1 = \text{corr}(U, V), \quad \rho_2 = \text{corr}(U^2 - U, V^2 - V),$$

where $(U, V) \sim C$. By using the following formula

$$\text{cov}(\nu(U), \xi(V)) = \int_{I^2} [C(u, v) - uv] d\nu(u) d\xi(v),$$

which provides the covariance between functions of U and V (Cuadras, 2002), we have

$$\begin{aligned}\rho_1 &= 12 \int_{I^2} C(u, v) dudv - 3, \\ \rho_2 &= 180 \int_{I^2} C(u, v)(4uv - 2u - 2v + 1) dudv - 5.\end{aligned}$$

Of course, ρ_1 is Spearman's rho correlation.

Expanding the AMH copula C_t , which plays the role of "true" copula,

$$\frac{uv}{1 - \theta(1-u)(1-v)} = uv \left[1 + \sum_{i=1}^{\infty} \theta^i (1-u)^i (1-v)^i \right],$$

we obtain

$$\rho_1 = 12 \sum_{i=1}^{\infty} \theta^i B(2, i+1)^2,$$

and

$$\rho_2 = 180 \sum_{i=1}^{\infty} \theta^i [4B(3, i+1)^2 - 4B(2, i+1)B(3, i+1) + B(2, i+1)^2],$$

where $B(\cdot, \cdot)$ is the beta function. All pairs (ρ_1, ρ_2) belong to \mathcal{R} and the AMH copula can be approximated by

$$C_2 = uv + \rho_1 3u(1-u)v(1-v) + \rho_2 5(2u^3 - 3u^2 + u)(2v^3 - 3v^2 + v).$$

Kendall's tau $\tau = 4 \int_{I^2} CdC - 1$ for copula C_2 is (Nelsen *et al.*, 1997):

$$\tau(C_2) = \frac{2}{3}\rho_1 + \frac{2}{15}\rho_1\rho_2.$$

See Nelsen (2006) for the exact expressions for τ and ρ_S in the AMH family.

If $M = \min\{u, v\}$ and $W = \max\{u + v - 1, 0\}$ are the Fréchet–Hoeffding upper and lower bounds, two measures of fit are

$$\begin{aligned}\eta_1 &= \max_{u, v \in I} |C_t(u, v) - C_2(u, v)|, \\ \eta_2 &= D(C_t, C_2)/D(M, W),\end{aligned}$$

where

$$D(C_t, C_2) = \int_{I^2} (C_t - C_2)^2 dudv,$$

which satisfies $D(C_t, C_2) < D(M, W) = 1/24$. Thus $0 < \eta_i < 1$, $i = 1, 2$.

TABLE 1. Canonical correlations and fit for the AMH copula.

θ	ρ_1	ρ_2	η_1	η_2	$\rho_S(AMH)$	$\tau(AMH)$	$\tau(C_2)$
-1	-0.2711	0.0217	0.0055	0.0002	-0.2710	-0.1817	-0.1815
-0.5	-0.1489	0.0080	0.0017	0.0000	-0.1489	-0.0995	-0.0995
.5	0.1924	0.0223	0.0032	0.0001	0.1924	0.1288	0.1286
1	0.4783	0.2323	0.0261	0.0029	0.4784	0.3333	0.3335

Table 1 reports a numerical illustration, showing that the fit is quite good, practically preserving Spearman's rho and Kendall's tau.

5.2. Clayton–Oakes copula. We consider the Clayton–Oakes copula (Nelsen, 2006):

$$C_t = [\max(u^{-\theta} + v^{-\theta} - 1, 0)]^{-1/\theta}, \quad -1 \leq \theta < \infty.$$

The correlations ρ_1 and ρ_2 have been obtained numerically. However, now the results can provide FGM approximations which are not copulas, i.e., the density $c(u, v)$ is negative for some values of u, v . Then we take $(\rho_1^*, \rho_2^*) \in \mathcal{R}$ with the smallest Euclidean distance to (ρ_1, ρ_2) . Thus the Clayton–Oakes copula can be approximated by

$$C_2 = uv + \rho_1^* 3u(1-u)v(1-v) + \rho_2^* 5(2u^3 - 3u^2 + u)(2v^3 - 3v^2 + v).$$

TABLE 2. Canonical correlations and fit for the Clayton–Oakes copula.

θ	ρ_1	ρ_2	ρ_1^*	ρ_2^*	η_1	η_2	$\lambda_L(CO)$
-1	-1	1	-0.5085	0.2950	0.0625	0.0060	-
-0.5	-0.4667	0.1997	-0.4665	0.1995	0.0263	0.0036	-
0.5	0.2950	0.1150	0.2950	0.1150	0.0162	0.0014	0.2500
1	0.4784	0.2337	0.4785	0.2335	0.0261	0.0029	0.5000
2	0.6822	0.4104	0.5495	0.2620	0.0349	0.0035	0.7071
5	0.8846	0.6809	0.5190	0.2875	0.0385	0.0030	0.8706
10	0.9582	0.8470	0.5190	0.2875	0.0387	0.0037	0.9330

The fit is acceptably good for intermediate (positive) values of the parameter θ , see Table 2. However, there are differences in the upper and lower tail dependence parameters λ_L, λ_U . In both families $\lambda_U = 0$, whereas $\lambda_L = 0$ for C_2 . The Clayton–Oakes copula has $\lambda_L = 2^{-1/\theta}$ if $\theta \geq 0$, but there is no mass in the lower-left corner if $\theta < 0$, so λ_L does not exist if the parameter is negative.

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DEPARTMENT OF STATISTICS, FACULTY OF BIOLOGY, UNIVERSITY OF BARCELONA,
BARCELONA, SPAIN

E-mail address: ccuadras@ub.edu

DEPARTMENT OF MATHEMATICS, FACULTY OF ECONOMICS, UNIVERSITY OF ANTIOQUIA,
MEDELLIN, COLOMBIA

E-mail address: wdiaz@udea.edu.co